

A LETTERS JOURNAL EXPLORING THE FRONTIERS OF PHYSICS



### LETTER

## Steering and focusing of fundamental shear horizontal guided waves in plates by using multiple-strip metasurfaces

To cite this article: Hao Qiu et al 2019 EPL 127 46004

View the article online for updates and enhancements.

### You may also like

- Manipulation of shear horizontal guided wave with arbitrary wave fronts by using metasurfaces Hao Qiu and Faxin Li
- <u>SH guided wave excitation by an apparent</u> <u>face-shear mode (d<sub>36</sub>) piezocomposite</u> <u>transducer: experiments and theory</u> Hongchen Miao, Lei Xu and Hao Zhang
- Excitation and manipulation of guided shear-horizontal plane wave using elastic metasurfaces
- Xi Cao, Chunlin Jia, Hongchen Miao et al.



# Steering and focusing of fundamental shear horizontal guided waves in plates by using multiple-strip metasurfaces

HAO QIU<sup>1,2</sup>, MINGTONG CHEN<sup>1</sup>, QIANG HUAN<sup>1</sup> and FAXIN LI<sup>1,2(a)</sup>

<sup>1</sup> LTCS and College of Engineering, Peking University - Beijing 100871, China
<sup>2</sup> Center for Applied Physics and Technology, Peking University - Beijing, China

received 9 June 2019; accepted in final form 27 August 2019 published online 1 October 2019

PACS 62.30.+d – Mechanical and elastic waves; vibrations PACS 77.65.-j – Piezoelectricity and electromechanical effects

**Abstract** – We proposed a metasurface consisting of multiple parallel strips with varied widths which can steer and focus the fundamental shear horizontal (SH<sub>0</sub>) wave in plates. Firstly, the group velocity of the quasi-SH<sub>0</sub> wave in different-width strips was calculated. Then, a multiple-strip metasurface was designed which was realized by through-thickness grooving in an aluminum plate. Finite-element simulations show that this metasurface can steer the SH<sub>0</sub> wave to a prescribed angle. Later, experiments were conducted in which thickness-shear ( $d_{15}$ ) piezoelectric strips were used as SH<sub>0</sub> wave exciters and receivers, and experimental results well reproduced the simulation results. Finally, another multiple-strip metasurface was designed to focus the SH<sub>0</sub> wave at a designated point. Both simulations and experiments were conducted, and results validated the proposed design. This work may provide guidance for the design of SH wave metasurfaces.

Copyright © EPLA, 2019

Metasurface, whose definition is in analogy to the metamaterials [1], is an artificial interface typically consisting of subwavelength structures and can steer or reflect the incident wave in an anomalous way [2]. The concept of metasurface was firstly raised in optics [3–8], then extended to acoustics [9–14], and was recently applied to elastic waves [15–23]. The working principle of all the metasurface is based on the Generalized Snell's Law (GSL) which indicates that when there exists a phase gradient at the interface, abnormal reflection and refraction may appear.

Unlike the optical or acoustical metasurfaces which are mostly for bulk waves, the elastic metasurfaces are typically for guided waves in finite-dimension solids. Yan *et al.* firstly proposed and numerically validated a metasurface (which they called surface bonded elastic metameterial) by bonding a two-dimensional planar array of small lead discs on an aluminum plate with silicone rubber to bend and focus low-frequency flexural  $A_0$  Lamb waves [21]. Zhu and Semperlotti shifted the phase of incoming wave by locally resonant toruslike tapers and realized anomalous refraction of symmetric (S<sub>0</sub>) or antisymmetric (A<sub>0</sub>) Lamb waves [24]. Liu *et al.* proposed a zigzag, 3-D printed metasurface which can shift and split a point source via forming phase gradient for the flexural Lamb wave  $(A_0)$  [15]. Lee *et al.* proposed a high transmission metasurface by mass-stiffness substructuring and successfully realized anomalous steering and focusing of the inplane symmetric Lamb wave  $(S_0)$  [18]. Cao *et al.* proposed a 3-D metasurface consisting of multiple slender pillars for full phase shift and high transmission, which can realize abnormal deflection and negative refraction of flexural Lamb wave  $(A_0)$  [17]. They also realized asymmetric flexural wave  $(A_0)$  transmission based on dual-layer elastic wave gradient metasurfaces [16]. Recently, Zhu et al. realized total internal reflection of the flexural Lamb wave  $(A_0)$  based on elastic metasurface consisting of three types of unit cells and applied it to structural vibration isolation [23]. Besides the above-mentioned metasurfaces for Lamb waves, metasurfaces for Rayleigh wave [20], Love wave [19] and bulk SH wave [25] were also proposed and simulated, but experimental validation is lacking.

In this letter, we proposed a metasurface for a guided fundamental shear horizontal  $(SH_0)$  wave which consists of several parallel strips with varied width and was fabricated by simply grooving on an aluminum plate. Firstly, simulation results show that the group velocity of the quasi-SH<sub>0</sub> wave in the strips increases steadily with the increasing strip width. Then, a multiple-strip metasurface was proposed with the strip width increasing monotonically. Simulations indicate that when a plane incident SH<sub>0</sub> wave

<sup>(</sup>a)E-mail: lifaxin@pku.edu.cn

passes through this metasurface, a phase gradient will be generated and the  $SH_0$  wave will be steered. Later, experiments were conducted on an aluminum plate using thickness-shear  $(d_{15})$  piezoelectric strips for  $SH_0$  wave excitation and reception, and the experimental results validated the proposed design. Finally, another multiplestrip metasurface was proposed to focus the  $SH_0$  wave, both simulations and experiments validated the proposed design.

According to the generalized Snell's law (GSL) [26], when there exists a phase gradient  $(d\phi/dy)$  of the incident wave at the interface between two media, the relationship between the refraction angle  $\theta_t$  and the incident angle  $\theta_i$ turns to be

$$\frac{\sin(\theta_t)}{\lambda_t} - \frac{\sin(\theta_i)}{\lambda_i} = \frac{1}{2\pi} \frac{\mathrm{d}\phi}{\mathrm{d}y},\tag{1}$$

where  $\lambda_i$  and  $\lambda_t$  are the wavelengths in the incident medium and the transmission medium, respectively.

Obviously, based on eq. (1), when an interface (metasurface) is introduced in a uniform medium which can cause the phase gradient of the incident wave, abnormal refraction will occur, *i.e.*, the incident wave will be steered in the uniform medium by the metasurface. In this work, the phase gradient is caused by the different group velocities in strips of different widths.

Based on the early work by Mindlin and Fox [27], Cegla calculated the dispersive curve of the quasi- $SH_0$  wave in an infinite-long, large aspect ratio rectangular strip and found that at high frequency above 200 kHz, the phase (group) velocity decreases (increases) with increasing strip width [28]. In this work, the working frequency of the proposed metasurface is set as 100 kHz and the aspect ratio of the rectangular strip cross-section cannot be very large, thus Cegla's results [28] cannot be employed here. We firstly simulated the guided wave propagation in 30 mm long, 2 mm thick aluminum strips with varied widths using COMSOL. In the simulation, a  $100 \,\mathrm{kHz}$  plane  $\mathrm{SH}_0$  wave was excited by a 100 mm long shearing line force on a large aluminum plate, then propagated into the strip and finally propagated to another large aluminum plate, as shown in the insert of fig. 1(a). Results show that the SH-type guided wave is dominant in the strip, see fig. S1 in the supplementary material Supplementarymaterial.pdf (SM), and here it is called quasi-SH<sub>0</sub> wave.

The group velocity of the quasi-SH<sub>0</sub> in strips with different frequency-width products was calculated and the results were plotted in fig. 1(a) in which the phase shift was also presented with the 200 kHz  $\cdot$  mm strip case as the reference.

As the group velocities of the quasi- $SH_0$  wave in a strip is dependent on the frequency-width product, the design of the metasurface (width of the strips) needs to be changed accordingly for different frequencies.

Considering the case of 100 kHz, it can be seen from fig. 1(a) that the group velocity of the quasi-SH<sub>0</sub> in the

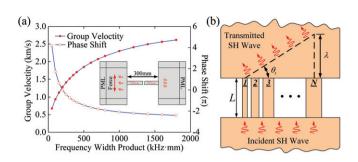


Fig. 1: (a) Simulated quasi- $SH_0$  wave propagation in 30 mm long, 2 mm thick aluminum strips with varied frequency-width products. Dots: group velocity; circles: phase shift referring to the 200 kHz.mm strip. (b) Illustration of the working principle of the multiple-strip metasurface.

strip increases steadily and nonlinearly with the increasing width. The phase shift referring to the  $2 \,\mathrm{mm}$  wide strip can vary from  $-2\pi$  to  $4\pi$  in the width range [0.5 mm,  $18 \,\mathrm{mm}$ ]. For practical design, the width range  $[2 \,\mathrm{mm}]$ , 8 mm] was selected and the phase shift can cover a full  $2\pi$ range. Here to design a metasurface for steering the  $SH_0$ wave by an angle  $\theta_t$ , multiple parallel 30 mm long strips were proposed with the strip width increasing steadily from  $2 \,\mathrm{mm}$  to  $8 \,\mathrm{mm}$ , as shown in fig. 1(b). The multiple parallel strips were fabricated by simply through-thickness grooving on a large aluminum plate. The gap between all adjacent strips was set to the same, and denoted by  $W_{gap}$ . The design process of the multiple-strip SH wave metasurface is as follows. As seen in fig. 1(b), the width of the left, 1st strip is  $W_1 = 2 \text{ mm}$ . The width of the 2nd strip,  $W_2$ , should satisfy the following implicit equation:

$$\frac{(W_1/2 + W_{gap} + W_2/2)}{\Delta\phi(W_2)} \tan(\theta_t) = \frac{\lambda_{SH}}{2\pi},$$
 (2)

where the function of  $\Delta \phi(W)$  can be obtained by curve fitting of the calculated phase shift in fig. 1(a).

 $W_2$  can be obtained by solving eq. (2) numerically. Similarly, the width of the *i*-th strip,  $W_i$ , can be obtained in sequence.

Simulations and experiments. - Finite-element (FEM) simulations were then conducted by using the commercial software COMSOL to predict the performances of the design metasurface. Here the steering angle of the  $SH_0$  wave is set to be  $18.5^\circ$ , and the designed metasurface consists of 17 strips with the width increasing from 2 mm to 8 mm. For the convenience of fabrication, the gap between adjacent strips was set to be 2 mm. The length of the metasurface is about 94 mm and the width of each strip is provided in the SM. The schematic of the FEM model was shown in fig. 2(a), in which the 100 kHz plane  $SH_0$  wave was excited by using 10 pieces of 20 mm long  $d_{15}$ piezoelectric strips bonded on the aluminum plate in one row with a gap of 4 mm. The PML (Perfectly Matched Layers) were used at all boundaries to remove the boundary reflections. The simulated in-plane Y displacement of

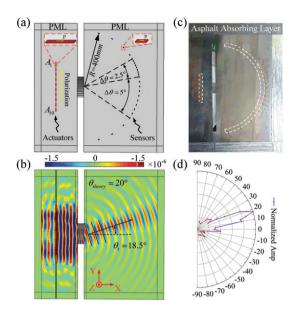


Fig. 2: Steering of the SH<sub>0</sub> wave in an aluminum plate using the proposed multiple-strip metasurface. (a) Schematic of finiteelement simulation and experimental testing; (b) photo of the experimental setup; (c) simulated in-plane displacement of the SH<sub>0</sub> wave; (d) amplitude of the steered SH<sub>0</sub> wave received by  $d_{15}$  piezoelectric strips.

the incident  $SH_0$  wave and that of the transmitted  $SH_0$ wave through the metasurface were shown in fig. 2(b). It can be seen that the designed metasurface can indeed steer the  $SH_0$  wave by an angle of 18.5°. The energy transmission rate by the designed metasurface is about 65%, which is fairly high. The energy loss is mainly caused by the relatively large gap (2 mm) between the adjacent strips, which will reflect a considerable amount of  $SH_0$  wave energy. If the gap were reduced, the energy transmission rate can be considerably increased.

To further validate the design of the metasurface, experiments were conducted on a  $1200 \times 1200 \times 2 \text{ mm}^3$  aluminum plate on which the multiple-strip metasurface was fabricated by simply through-thickness grooving. By applying asphalt on the upper and lower surfaces of the four edges of the aluminium plate (15 mm width asphalt sticker is used in this experiment), the SH wave can be effectively absorbed and the effect of "absorbing layer" can be realized. The  $SH_0$  wave was excited by using 10 pieces of 20 mm long  $d_{15}$  piezoelectric strips, the same as those in the simulation. The steered  $SH_0$  wave was received by 8 mm long  $d_{15}$ piezoelectric strips distributed non-uniformly on a 400 mm radius arc with the middle point of the metasurface as the circle center, as shown in fig. 2(a) and (c), where the best angle resolution in measurement is  $2.5^{\circ}$ . Figure 2(d) shows the measured amplitude of the steered  $SH_0$  wave distributed along the arc sensor array. It can be seen that, as expected, the  $SH_0$  wave was steered by the designed metasurface. The maximum amplitude appears at  $17.5^{\circ}$ . Considering the measurement resolution of  $2.5^{\circ}$ , the measured steering angle should be  $16.25^{\circ}-18.75^{\circ}$ , which is

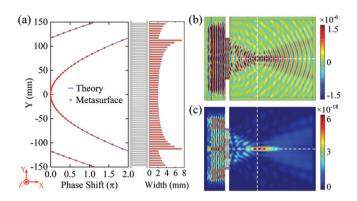


Fig. 3: Simulated focusing of the  $SH_0$  wave using a multiplestrip metasurface. (a) Phase shift corresponding to the strip width; (b) simulated in-plane Y displacement of the  $SH_0$  wave showing the focusing point; (c) simulated  $SH_0$  wave energy highlighting the focusing point.

consistent with the designed steering angle of  $18.5^{\circ}$ . This confirmed the validity of the designed metasurface in the steering of the SH<sub>0</sub> wave.

Based on the above  $SH_0$  wave steering results, we further proposed a multiple-strip metasurface for the SH wave focusing with the focal length of 200 mm. The designed metasurface for a 100 kHz  $SH_0$  wave consists of 61 strips and is illustrated in fig. 3(a) in which the calculated and theoretical phase shift was also presented. The length of the focusing metasurface is about 300 mm and the detailed width info is provided in the SM. Also for the convenience of fabrication, the gap between adjacent strips is set to 2 mm.

Finite-element simulations were firstly conducted to predict the performance of the designed metasurface in SH wave focusing. In the simulation, the 100 kHz SH<sub>0</sub> wave was excited by using 20 pieces of 20 mm long  $d_{15}$  piezoelectric strips in one row bonded on the aluminum plate. Figure 3(b) shows the simulated in-plane Y displacement of the incident SH<sub>0</sub> wave and that of the transmitted SH<sub>0</sub> wave via the metasurface. It can be seen that the SH<sub>0</sub> wave was focused by the designed metasurface with a focal length of 200 mm. After propagating across the focal point, the SH<sub>0</sub> wave turns to be divergent. The simulated energy distribution of the SH<sub>0</sub> wave was shown in fig. 3(c), from which it can be seen more clearly that the wave energy was focused at the designed focal point.

Experiments were then conducted to check the performance of the designed focusing metasurface. In the same way as that in the simulation, the incident 100 kHz SH<sub>0</sub> wave was excited by 20 pieces of 20 mm long  $d_{15}$  piezoelectric strips. To reduce the effect of the bonded sensors on the received SH wave signals, two sets of  $d_{15}$  strip sensor arrays were bonded in sequence around the designed focal point, *i.e.*, along the X-direction (fig. 4(a)) and Y-direction (fig. 4(c)) of the designed focal point, respectively. Figure 4(b) and fig. 4(d) show the measured SH<sub>0</sub> wave energies along the X-direction and Y-direction

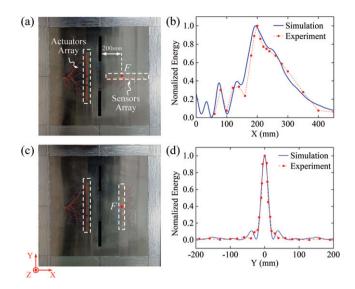


Fig. 4: Experimental verification of the  $SH_0$  wave focusing using a metasurface. Left: layout of the experimental setup showing the different sensor arrays around the designed focal point, (a) along the X-direction, (c) around the Y-direction; Right: Measured and simulated  $SH_0$  wave signals in different sensor arrays, (b) along the X-direction, (d) along the Y-direction.

of the focal point, respectively. It can be seen that the experimental results are consistent with the simulation results for both cases, *i.e.*, the  $SH_0$  wave energy reaches its maxima at the designed focal point along both directions. From fig. 4, it can be concluded that the designed metasurface can effectively focus the SH wave at the designed focal point.

It should be noted that in above simulations and experiments, the excitation waves are always continuous waves. Actually, this metasurface can also work under pulse excitations. To show this, we have simulated the performance of the proposed metasurface under pulse excitation using the semi-analytical finite-element method (SAFE) [29,30] and the results were presented in fig. S3 and fig. S4 in the SM. However, due to the dispersion of the quasi-SH<sub>0</sub> wave in the strips, there exist a quasi-SH wave with different group velocities in the same strip under pulse excitation signals with a certain frequency bandwidth, which will affect the performance of the metasurface and bring difficulties to the design. How to make this metasurface work better under pulse excitations requires further investigations.

In summary, we proposed a design of metasurface for the  $SH_0$  wave steering and focusing in plates based on multiple parallel strips. The metasurface is based on the generalized Snell's law and the rule that the group velocity of quasi-SH<sub>0</sub> in strips increases steadily with the strip width. Both simulations and experimental measurement validated the designed metasurfaces in  $SH_0$  wave steering and focusing. The proposed multiple-strip metasurface is easy to fabricate and the transmission rate can be fairly high. This work provides a convenient solution for manipulating the SH wave in plate-like structures.

\* \* \*

Financial support from the National Natural Science Foundation of China under Grant No. 11672003 is greatly acknowledged.

#### REFERENCES

- SMITH D. R., PENDRY J. B. and WILTSHIRE M. C. K., Science, **305** (2004) 788.
- [2] YU N. F. and CAPASSO F., Nat. Mater., 13 (2014) 139.
- [3] FANG N. et al., Science, **308** (2005) 534.
- [4] LIN J. et al., Science, **340** (2013) 331.
- [5] NI X. J. et al., Science, **335** (2012) 427.
- [6] NOVOTNY L. and VAN HULST N., Nat. Photon., 5 (2011) 83.
- [7] SCHURIG D. et al., Science, **314** (2006) 977.
- [8] SHELBY R. A., SMITH D. R. and SCHULTZ S., Science, 292 (2001) 77.
- [9] CHEN J. et al., Appl. Phys. Lett., 114 (2019) 104101.
- [10] JIANG X. et al., Phys. Rev. Lett., **117** (2016) 034301.
- [11] LI Y. et al., Phys. Rev. Appl., 4 (2015) 024003.
- [12] LI Y. et al., Phys. Rev. Appl., 101 (2012) 233508.
- [13] MA G. et al., Nat. Mater., **13** (2014) 873.
- [14] XIE Y. et al., Nat. Commun., 5 (2014) 5553.
- [15] LIU Y. Q. et al., Phys. Rev. Lett., **119** (2017) 034301.
- [16] CAO L. Y. et al., Appl. Phys. Lett., **113** (2018) 183506.
- [17] CAO L. Y. et al., Smart Mater. Struct., 27 (2018) 075051.
- [18] LEE H. et al., J. Mech. Phys. Solids, **112** (2018) 577.
- [19] PALERMO A. and MARZANI A., Sci. Rep., 8 (2018) 7234.
- [20] XU Y. L., YANG Z. C. and CAO L. Y., J. Phys. D: Appl. Phys., 51 (2018) 175106.
- [21] YAN X. et al., Appl. Phys. Lett., 103 (2013) 121901.
- [22] ZHU H. F. and SEMPERLOTTI F., Phys. Rev. Lett., 117 (2016) 034302.
- [23] ZHU H. F., WALSH T. F. and SEMPERLOTTI F., Appl. Phys. Lett., 113 (2018) 221903.
- [24] ZHU H. and SEMPERLOTTI F., Phys. Rev. Lett., 117 (2016) 034302.
- [25] CAO L. Y., YANG Z. C. and XU Y. L., J. Sound Vib., 418 (2018) 1.
- [26] YU N. F. et al., Science, **334** (2011) 333.
- [27] MINDLIN R. D. and FOX E. A., J. Appl. Mech., 27 (1960) 152.
- [28] CEGLA F. B., J. Acoust. Soc. Am., 123 (2008) 4218.
- [29] ZUO P., YU X. D. and FAN Z., NDT & E Int., 90 (2017) 11.
- [30] ZUO P. and FAN Z., J. Sound Vib., 406 (2017) 181.